

Direct images of secondary classes and renormalization

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- I. Chiral anomalies and spectral triples (not new)
- II. Conformal renormalization: example (some news)
- III. Direct images of secondary classes (new)

I. Chiral anomalies and spectral triples

- Let (\mathcal{C}, H, D) be a spectral triple of even degree and finite summability.

\mathcal{C} = Fréchet algebra represented in Hilbert space

$$H = \begin{pmatrix} H_+ \\ H_- \end{pmatrix}, \quad D = \begin{pmatrix} 0 & D_- \\ D_+ & 0 \end{pmatrix}$$

- Chiral gauge theory: classically

matter fields $\psi \in H_+$, $\bar{\psi} \in (H_-)^*$

gauge potential $A = \sum_{i,j} c_i^0 [D, c_j^1]$

The classical dynamics of the matter fields evolving in the potential A is governed by the action functional

$$\Gamma^{(0)}(\psi, \bar{\psi}, A_+) = \langle \bar{\psi}, (D_+ + A_+)\psi \rangle$$

Suppose D invertible. The quantum corrections to the classical action (A remains fixed) is the formal power series

$$\Gamma^{(1)}(A_+) = \text{Tr}(D_+^{-1} A_+) - \frac{1}{2} \text{Tr}((D_+^{-1} A_+)^2) + \frac{1}{3} \text{Tr}((D_+^{-1} A_+)^3) \dots$$

$$= \text{loop} - \frac{1}{2} \text{bubble} + \frac{1}{3} \text{triangle} - \frac{1}{4} \text{square} + \frac{1}{5} \text{pentagon} + \dots$$

Low-order graphs have to be renormalized (heat operator, zeta functions, Dim Reg, etc ...)

- Gauge symmetry: Let $u \in \mathcal{C} \hat{\otimes}_\pi C^\infty(S^1)$ be a smooth loop of invertibles (unitaries) in \mathcal{C} .

Gauge transformations of $(\psi_0, \bar{\psi}_0, A_0)$:

$$\begin{aligned} \psi &= u_+^{-1} \psi_0, & \bar{\psi} &= \bar{\psi}_0 u_- \\ A &= u^{-1} [D, u] + u^{-1} A_0 u \end{aligned}$$

Fact: the classical action $\Gamma^{(0)}(\psi, \bar{\psi}, A_+)$ is constant over S^1 but not the quantum corrections $\Gamma^{(1)}(A_+)$.

\Rightarrow **anomaly**

The differential of the formal power series $\Gamma^{(1)}(A)$ is a *finite sum*, polynomial in A , linear in $\omega = u^{-1} s u$, and **local**. It defines a de Rham cohomology class of the circle

$$s\Gamma^{(1)}(A_+) = \Delta(\omega, A_+) \in H^1(S^1)$$

Theorem (P. 2005):

1) Let $\mathcal{A} = \mathcal{C} \hat{\otimes}_\pi C^\infty(S^1)$ and $\mathcal{B} = C^\infty(S^1)$. For any renormalization of $\Gamma^{(1)}(A)$, one has the commutative diagram

$$\begin{array}{ccc} K_1^{\text{top}}(\mathcal{A}) & \longrightarrow & HP_1(\mathcal{A}) \\ \text{index}(D) \downarrow & \searrow \Delta & \downarrow \text{ch}(D) \\ K_1^{\text{top}}(\mathcal{B}) & \longrightarrow & HP_1(\mathcal{B}) \\ \parallel & & \parallel \\ \mathbb{Z} & & \mathbb{C} \end{array}$$

2) If (\mathcal{C}, H, D) is *regular* and $\Gamma^{(1)}(A_+)$ is renormalized by zeta-function, the anomaly is a finite sum of residues:

$$\Delta(\omega, A_+) = \text{Res}_{z=0} \frac{1}{z} \text{Tr}(\gamma\omega |D|^{-2z}) + \sum_{n \geq 1, k \geq 0} (-1)^{n+k} c(k) \times$$

$$\text{Res}_{z=0} \left(\frac{\Gamma(z+n+k)}{z\Gamma(z)} \text{Tr}((q\omega A^{(k_1)} D A^{(k_2)} \dots D A^{(k_n)} |D|^{-2(z+n+k)}) \right)$$

$k = (k_1, \dots, k_n)$ is a multi-index

$$q\omega = \frac{1+\gamma}{2} [\omega, D]$$

$A^{(k_i)}$ = k_i -th commutator of D^2 on A

$$c(k)^{-1} = (k_1! \dots k_n!) (k_1+1)(k_1+k_2+2) \dots (k_1+\dots+k_n+n)$$

II. Conformal renormalization: simple example

- Smooth étale groupoid $G \ltimes \Sigma$:

$\Sigma = \mathbb{C}$ space of units, G pseudogroup of conformal diffeomorphisms $g : \text{Dom}(g) \subset \Sigma \rightarrow \text{Ran}(g) \subset \Sigma$

- Either g has finitely many fixed points $z_0 \in \text{Dom}(g)$, with $g'(z_0) \neq 1$,

- Or $g = \text{id}$ on $\text{Dom}(g)$.

Crossed product algebra $\mathcal{C} = C_c^\infty(\Sigma) \rtimes G$ linearly generated by symbols fU_g^* , $f \in C_c^\infty(\text{Dom}(g))$, and product

$$(f_1 U_{g_1}^*)(f_2 U_{g_2}^*) = f_1 (f_2)^{g_1} U_{g_2 g_1}^*$$

- Dolbeault operator $D_+ = \bar{\partial} = d\bar{z}\partial_{\bar{z}} : C^\infty(\Sigma) \rightarrow \Omega^{0,1}(\Sigma)$.

(\mathcal{C}, D_+) is not a spectral triple (no D_-), but D_+ is *conformally invariant*.


- Conformal field theory: chiral fields $\psi \in C^\infty(\Sigma)$, $\bar{\psi} \in \Omega_c^{1,0}(\Sigma)$ and gauge potential $A_+ = \sum_{i,j} c_i^0 [D_+, c_j^1]$

$$\Gamma^{(0)}(\psi, \bar{\psi}, A_+) = \int_{\Sigma} \bar{\psi} (D_+ + A_+) \psi$$

$$\text{Propagator: } D_+^{-1}(z, w) = \frac{1}{\pi(z-w)}$$

$$\Gamma^{(1)}(A_+) = \text{diagram 1} - \frac{1}{2} \text{diagram 2} + \frac{1}{3} \text{diagram 3} - \frac{1}{4} \text{diagram 4} + \frac{1}{5} \text{diagram 5} + \dots$$

Only the first two graphs need renormalization (Operator Product Expansion):

 integrates the kernel $k(z) = \frac{1}{g(z) - z}$ for some $g \in G$


- If g has no fixed point, then $k(z)$ is a smooth function (OK)

- If g has isolated fixed point z_0 , with $g'(z_0) \neq 1$. Then k is singular at z_0 with

$$k(z) \sim \frac{1}{(g'(z_0) - 1)} \frac{1}{(z - z_0)} \quad (\text{OK})$$

- If $g = \text{id}$ on $\text{Dom}(g)$, then $k(z)$ is not defined.

$$k(z) \rightsquigarrow 0$$

 integrates the kernel $k(z, w) = \frac{1}{(g(z) - w)(h(w) - z)}$

for some $g, h \in G$.

- If $gh \in G$ has only isolated fixed points OK.

- If $gh = \text{id}$ on $\text{Dom}(gh)$ then around $g(z) \sim w$ the function $k(z, w)$ exhibits a singularity of the kind

$$\frac{1}{(g(z) - w)^2} \rightsquigarrow \partial_w \frac{1}{(g(z) - w)}$$

- The chiral anomaly $\Delta(\omega, A_+) = s\Gamma^{(1)}(A_+)$ may be decomposed as a sum

$$\Delta(\omega, A_+) = \sum_{g \in G} \Delta(\omega, A_+)[g]$$

- If g has isolated fixed points then $(\omega = \sum_g \omega[g]U_g^*)$

$$\Delta(\omega, A_+)[g] = \sum_{z_0} \frac{1}{1 - g'(z_0)} \omega(z_0)[g]$$

- If $g = \text{id}$ on $\text{Dom}(g)$ then

$$\Delta(\omega, A_+)[g] = \frac{1}{2\pi i} \int_{\Sigma} (\partial\omega - \frac{1}{2}\delta\omega)A_+ [g]$$

with $\partial = dz\partial_z$ and $\delta = [\partial, \sigma]$, where σ is the generator of the modular automorphism group of $C_c^\infty(\Sigma) \rtimes G$:

$$\sigma(fU_h^*) = f \ln |h'|^2 U_h^*$$

Remark: With $\mathcal{A} = C^\infty(\Sigma \times S^1) \rtimes G$ and $\mathcal{B} = C^\infty(S^1)$ there exists a cyclic cocycle $\varphi : HP_1(\mathcal{A}) \rightarrow HP_1(\mathcal{B}) = \mathbb{C}$ such that

$$\begin{array}{ccc} K_1^{\text{top}}(\mathcal{A}) & \longrightarrow & HP_1(\mathcal{A}) \\ & \searrow \Delta & \downarrow \varphi \\ & & HP_1(\mathcal{B}) \end{array}$$

III. Direct images of secondary classes

All algebras are Fréchet, with topology generated by sub-multiplicative seminorms: $q(ab) \leq q(a)q(b)$ (needed for Bott periodicity).

Examples: Banach algebras, $C^\infty(M)$, smooth compact operators (matrices with rapid decay over \mathbb{C}), etc...

Primary and Secondary classes (Connes, Karoubi): For any such algebra \mathcal{A} there exists a long exact sequence

$$\dots K_{n+1}^{\text{top}}(\mathcal{A}) \rightarrow HC_{n-1}(\mathcal{A}) \xrightarrow{\delta} MK_n(\mathcal{A}) \rightarrow K_n^{\text{top}}(\mathcal{A}) \rightarrow HC_{n-2}(\mathcal{A}) \dots$$

For $\mathcal{A} = C^\infty(M)$ the multiplicative K -theory $MK_n(\mathcal{A})$ has close connections with (smooth) Deligne cohomology \Rightarrow geometric invariants.

Quasihomomorphisms (Cuntz): A p -summable quasihomomorphism from \mathcal{A} to \mathcal{B} is a pair of homomorphisms

$$\rho_\pm : \mathcal{A} \rightrightarrows \mathcal{E} \triangleright \mathcal{I} \hat{\otimes}_\pi \mathcal{B}$$

from \mathcal{A} to an algebra \mathcal{E} , with $\mathcal{I} \hat{\otimes}_\pi \mathcal{B}$ a (closed) two-sided ideal in \mathcal{E} , and \mathcal{I} a p -summable algebra (typically $\mathcal{I} = \mathcal{L}^p(H)$). For any $a \in \mathcal{A}$,

$$\rho_+(a) - \rho_-(a) \in \mathcal{I} \hat{\otimes}_\pi \mathcal{B}$$

Theorem (P. 2006):

1) Let $\rho : \mathcal{A} \rightrightarrows \mathcal{E} \triangleright \mathcal{I} \hat{\otimes}_\pi \mathcal{B}$ be a $(p + 1)$ -summable quasihomomorphism, with p even integer. Suppose that $K_n^{\text{top}}(\mathcal{I} \hat{\otimes}_\pi \mathcal{B}) \cong K_n^{\text{top}}(\mathcal{B})$ and ρ verifies adequate properties (existence of (pro)-nilpotent extensions etc...). Then one has a transformation of degree $-p$

$$\begin{array}{ccccccc}
 \dots & K_{n+1}^{\text{top}}(\mathcal{A}) & \longrightarrow & HC_{n-1}(\mathcal{A}) & \longrightarrow & MK_n(\mathcal{A}) & \longrightarrow & K_n^{\text{top}}(\mathcal{A}) & \dots \\
 & \downarrow \rho! & & \downarrow \text{ch}(\rho) & & \downarrow & & \downarrow \rho! & \\
 \dots & K_{n+1-p}^{\text{top}}(\mathcal{B}) & \longrightarrow & HC_{n-1-p}(\mathcal{B}) & \longrightarrow & MK_{n-p}(\mathcal{B}) & \longrightarrow & K_{n-p}^{\text{top}}(\mathcal{B}) & \dots
 \end{array}$$

2) The bivariant Chern character $\text{ch}(\rho) \in HC^p(\mathcal{A}, \mathcal{B})$ is the boundary of a renormalized eta-cochain

$$\text{ch}(\rho) = \partial \eta_R$$

and is given by a *local formula*. The diagonal map

$$\begin{array}{ccc}
 K_1^{\text{top}}(\mathcal{A}) & \xrightarrow{\text{ch}_{p+1}} & HC_{p+1}(\mathcal{A}) \\
 \downarrow & \searrow \Delta & \downarrow \text{ch}(\rho) \\
 K_1^{\text{top}}(\mathcal{B}) & \longrightarrow & HC_1(\mathcal{B})
 \end{array}$$

is an exact generalization of the chiral anomaly (morally, $\Gamma^{(1)} = \eta_R \circ \text{ch}_{p+1}$).